

PLANE WAVES IN A CONDUCTING MEDIUM

II. GOOD CONDUCTING CASE

If $\varepsilon'' \gg \varepsilon'$ or $\left(\frac{\sigma}{\omega\varepsilon} \gg 1\right)$, conduction current dominates considerably over the displacement current. Then,

$$\alpha = \omega \sqrt{\frac{\varepsilon\mu}{2}} \left[\sqrt{1 + \left(\frac{\sigma}{\omega\varepsilon}\right)^2} - 1 \right]^{1/2} \approx \omega \sqrt{\frac{\varepsilon\mu}{2}} \left[\left(\frac{\sigma}{\omega\varepsilon}\right) - 1 \right]^{1/2}$$

$$\alpha \approx \omega \sqrt{\frac{\varepsilon\mu}{2}} \sqrt{\frac{\sigma}{\omega\varepsilon}} = \sqrt{\frac{\omega\mu\sigma}{2}} = \sqrt{\frac{2\pi f \mu\sigma}{2}} = \sqrt{\pi f \mu\sigma}$$

Then, $\alpha \approx \sqrt{\pi f \mu\sigma}$, α is a function of frequency.

$$\beta = \omega \sqrt{\frac{\varepsilon\mu}{2}} \left[\sqrt{1 + \left(\frac{\sigma}{\omega\varepsilon}\right)^2} + 1 \right]^{1/2} \approx \omega \sqrt{\frac{\varepsilon\mu}{2}} \left[\frac{\sigma}{\omega\varepsilon} + 1 \right]^{1/2}$$

$$\beta = \omega \sqrt{\frac{\varepsilon\mu}{2}} \sqrt{\frac{\sigma}{\omega\varepsilon}} = \sqrt{\pi f \mu\sigma}$$

$\alpha = \beta = \sqrt{\pi f \mu\sigma}$. Then,

$$\gamma = \alpha + j\beta = \sqrt{\pi f \mu\sigma} (1 + j)$$

Intrinsic Impedance:

$$\eta_c = \sqrt{\frac{\mu}{\varepsilon_c}} = \sqrt{\frac{\mu}{\varepsilon - j\frac{\sigma}{\omega}}} = \sqrt{\frac{\mu}{\varepsilon}} \left(1 - j\frac{\sigma}{\omega\varepsilon}\right)^{-1/2}$$

$$\eta_c \approx \sqrt{\frac{\mu}{\varepsilon}} \left(-j\frac{\sigma}{\omega\varepsilon}\right)^{-1/2} = \sqrt{\frac{\mu}{\varepsilon}} \frac{1}{\sqrt{-j\frac{\sigma}{\omega\varepsilon}}} = \sqrt{\frac{\mu\omega}{\sigma}} \sqrt{j} = \sqrt{\frac{\mu\omega}{\sigma}} \frac{(1+j)}{\sqrt{2}}$$

$$\eta_c \approx (1+j) \sqrt{\frac{\omega\mu}{2\sigma}}$$

$$\sqrt{\frac{w\mu}{2\sigma}} = \sqrt{\frac{w\mu\sigma}{2\sigma^2}} = \sqrt{\frac{w\mu\sigma}{2}} \frac{1}{\sigma} = \frac{\alpha}{\sigma}$$

$$\text{So, } \eta_c = (1+j) \frac{\alpha}{\sigma} = \frac{\alpha\sqrt{2}}{\sigma} \angle 45^\circ$$

The results show that \bar{H} lags \bar{E} by 45° in time.

Phase Velocity

$$u_p = \frac{w}{\beta} = \frac{w}{\sqrt{\frac{w\mu\sigma}{2}}} = \sqrt{\frac{2w}{\mu\sigma}}$$

The wavelength,

$$\lambda = \frac{2\pi}{\beta} = \frac{2\pi}{\sqrt{\pi f \mu \sigma}} = 2\sqrt{\frac{\pi}{f \mu \sigma}}$$

$$\lambda = \frac{2\pi}{\beta} = \frac{2\pi}{\alpha} = 2\pi \left(\frac{1}{\alpha} \right) = 2\pi\delta$$

Example:

Consider copper for which $\sigma = 5.80 \times 10^7$ (S/m) $\varepsilon = \varepsilon_0$, $\mu = \mu_0$. Take $f = 3\text{MHz}$.

Solution:

Let us first find the wavelength of a wave in free space at this frequency:

$$\lambda_0 = \frac{c}{f} = \frac{3 \times 10^8}{3 \times 10^6} = 100 \text{ meters}$$

Calculate the loss tangent:

$$\frac{\sigma}{\omega\varepsilon} = \frac{5.8 \times 10^7}{2\pi \times 3 \times 10^6 \times \frac{1}{36\pi} \times 10^{-9}} = 3.48 \times 10^{11} \gg 1$$

Copper behaves like a good conductor at $f = 3\text{MHz}$.

The wavelength in copper will then be:

$$\lambda = 2 \left(\frac{\pi}{f \mu \sigma} \right)^{1/2} = 2 \left(\frac{\pi}{3 \times 10^6 \times 4\pi \times 10^{-7} \times 5.8 \times 10^7} \right)^{1/2} = 0.239 \times 10^{-3} \text{ m} .$$

or

$$\lambda = 0.239 \text{ mm} \text{ (compare with 100 meters in air!)}$$

$$\alpha = \sqrt{\pi f \mu \sigma} = \left(\pi \times 3 \times 10^6 \times 4\pi \times 10^{-7} \times 5.8 \times 10^7 \right)^{1/2} = 2.62 \times 10^4 \text{ (Np/m)}$$

The depth of penetration:

$$\delta = \frac{1}{\alpha} = 0.38 \times 10^{-4} \text{ m} = 0.038 \text{ mm} .$$

$$\delta \text{ at } f = 10 \text{ GHz is } 0.66 \mu\text{m} .$$

$$\text{The frequency at which } \sigma = \omega \epsilon \text{ is } f = \frac{5.8 \times 10^7}{2\pi \epsilon_0} = 1.04 \times 10^{18} \text{ Hz}$$

Thus at frequencies of even several gigahertz, copper behaves like an excellent conductor. The material is actually a good conductor well into the ultraviolet region. The skin-depth at optical frequencies ($\omega \sim 10^{15}$) is roughly 10^{-8} m , which explains why metals are opaque.

Example:

Consider sea-water, for which $\sigma = 4 \text{ mhos/m}$, $\epsilon = 80\epsilon_0$, $\mu = \mu_0$. Take $f = 25 \text{ kHz}$.

Solution:

$$\frac{\sigma}{\omega \epsilon} = \frac{4}{2\pi \times 25 \times 10^3 \times 80 \times \frac{1}{36\pi} \times 10^{-9}} = 36000 \gg 1$$

Sea-water is good conductor at $f = 25 \text{ kHz}$.

$$\lambda_{\text{air}} = \frac{3 \times 10^8}{25 \times 10^3} = 12 \text{ km}$$

$$\lambda_{\text{sea-water}} = 2 \left(\frac{\pi}{f \mu \sigma} \right)^{1/2} = 2 \left(\frac{\pi}{25 \times 10^3 \times 4\pi \times 10^{-7} \times 4} \right)^{1/2} = \frac{1}{4\pi \times 10^{-2}} \left(\frac{1}{f_{\text{kHz}}} \right)^{1/2} \text{ or}$$

$$\delta = \frac{8}{\sqrt{f_{kHz}}} \text{ (meter)}.$$

f (kHz)	δ (m)
10	2.531
15	2.067
20	1.789
25	1.6
30	1.462
35	1.353
100	0.8