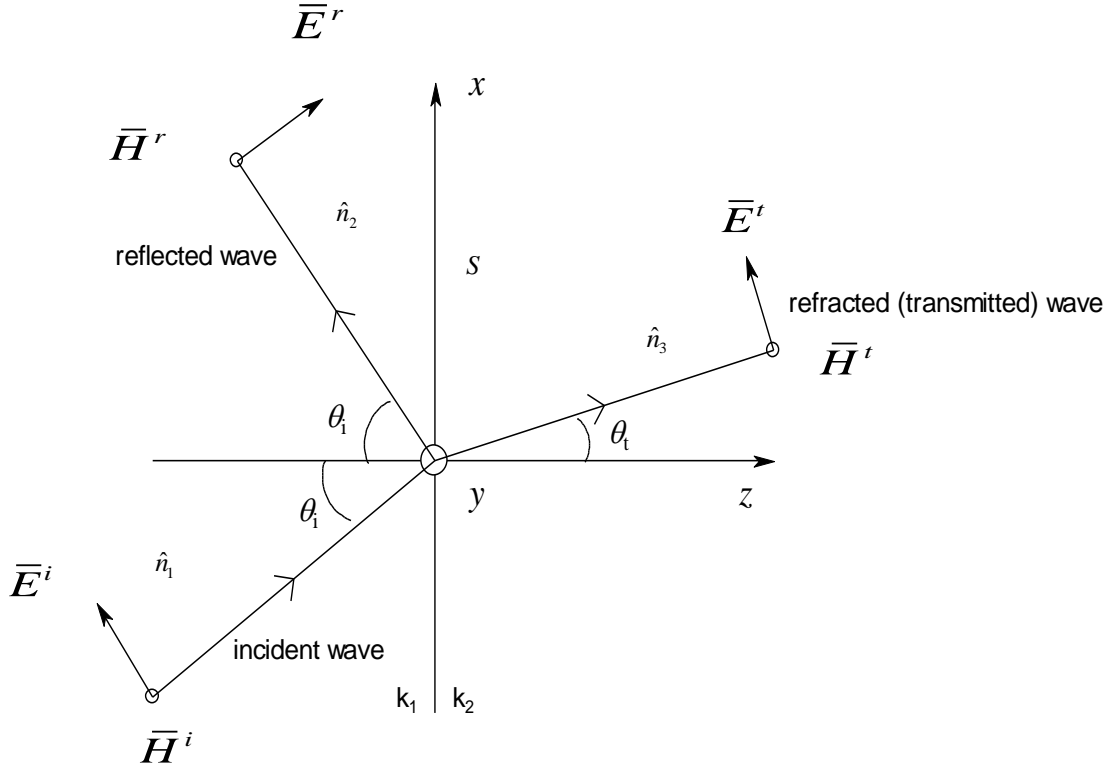


## PARALLEL POLARIZATION



$$\bar{H}^i = H_1 \hat{a}_y e^{-jk_1 \hat{n}_1 \cdot \bar{r}} \quad \bar{E}^i = \eta_1 \bar{H}^i \times \hat{n}_1$$

$$\bar{H}^r = -H_2 \hat{a}_y e^{-jk_1 \hat{n}_2 \cdot \bar{r}} \quad \bar{E}^r = \eta_1 \bar{H}^r \times \hat{n}_2$$

$$\bar{H}^t = H_3 \hat{a}_y e^{-jk_2 \hat{n}_3 \cdot \bar{r}} \quad \bar{E}^t = \eta_2 \bar{H}^t \times \hat{n}_3$$

$$\bar{E}^i = \eta_1 H_1 (\cos \theta_i \hat{a}_x - \sin \theta_i \hat{a}_z) e^{-jk_1 \hat{n}_1 \cdot \bar{r}}$$

$$\bar{E}^r = \eta_1 H_2 (\cos \theta_i \hat{a}_x + \sin \theta_i \hat{a}_z) e^{-jk_1 \hat{n}_2 \cdot \bar{r}}$$

$$\bar{E}^t = \eta_2 H_3 (\cos \theta_t \hat{a}_x - \sin \theta_t \hat{a}_z) e^{-jk_2 \hat{n}_3 \cdot \bar{r}}$$

Let,

$$E_1 = \eta_1 H_1, \quad E_2 = \eta_1 H_2, \quad E_3 = \eta_2 H_3$$

Boundary Conditions (after phase- matching) yield:

$$i) \quad (\bar{H}^i + \bar{H}^r)_{\tan} = (\bar{H}^t)_{\tan} \quad \text{at } z = 0 \quad \text{gives,}$$

$$H_1 - H_2 = H_3 \quad \text{or,}$$

$$\frac{1}{\eta_1}(E_1 - E_2) = \frac{1}{\eta_2}E_3 \quad \dots\dots\dots(1)$$

$$ii) \quad (\bar{E}^i + \bar{E}^r)_{\tan} = E_x^i + E_x^r = (\bar{E}^t)_{\tan} = E_x^t \quad \text{at } z = 0 \quad \text{gives}$$

$$(E_1 + E_2) \cos \theta_i = E_3 \cos \theta_t \quad \dots\dots\dots(2)$$

Define:

$$\frac{E_2}{E_1} = \Gamma_{\parallel}, \quad \frac{E_3}{E_1} = T_{\parallel} \quad \text{From (1) and (2) we get,}$$

## Fresnel Formulas for Parallel Polarization

$$\Gamma_{\parallel} = \frac{\eta_2 \cos \theta_t - \eta_1 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i}$$

$$T_{\parallel} = \frac{2\eta_2 \cos \theta_i}{\eta_2 \cos \theta_t + \eta_1 \cos \theta_i}$$

$\Gamma_{\parallel}$  is called the reflection coefficient for the parallel polarization.

$T_{\parallel}$  is called the transmission coefficient for the parallel polarization.

We see that,

$$1 + \Gamma_{\parallel} = \left( \frac{\cos \theta_t}{\cos \theta_i} \right) T_{\parallel}$$

Fresnel formulas are valid for  $\epsilon_1$  and  $\epsilon_2$  being complex (the two media are lossy). They are also valid for all frequencies.

## TOTAL TRANSMISSION-BREWSTER ANGLE

Question: For a given set of material parameters of the two media forming an interface, is there an incidence angle that allows no reflection, i.e.  $\Gamma = 0$ ?

a) Perpendicular Polarization

We will set  $\Gamma_{\perp} = 0$  and see under what conditions can this be possible.

$$\Gamma_{\perp} = \frac{\sqrt{\frac{\mu_2}{\epsilon_2}} \cos \theta_i - \sqrt{\frac{\mu_1}{\epsilon_1}} \cos \theta_t}{\sqrt{\frac{\mu_2}{\epsilon_2}} \cos \theta_i + \sqrt{\frac{\mu_1}{\epsilon_1}} \cos \theta_t} = 0$$

Or,

$$\cos \theta_i = \sqrt{\frac{\mu_1 \epsilon_2}{\mu_2 \epsilon_1}} \cos \theta_t \dots\dots\dots(1)$$

We also have from the Snell's Law:

$$\sqrt{\epsilon_1 \mu_1} \sin \theta_i = \sqrt{\epsilon_2 \mu_2} \sin \theta_t \dots\dots\dots(2)$$

We can have write (1) as:

$$(1 - \sin^2 \theta_i) = \frac{\mu_1 \epsilon_2}{\mu_2 \epsilon_1} (1 - \sin^2 \theta_t) \dots\dots\dots(3)$$

Let us now replace  $\sin \theta_t$  on the RHS of (3) with its equivalent from (2):

$$(1 - \sin^2 \theta_i) = \frac{\mu_1 \epsilon_2}{\mu_2 \epsilon_1} \left( 1 - \frac{\mu_1 \epsilon_1}{\mu_2 \epsilon_2} \sin^2 \theta_i \right)$$

$$\sin \theta_i = \left( \frac{\frac{\varepsilon_2 - \mu_2}{\varepsilon_1 \mu_1}}{\frac{\mu_1 - \mu_2}{\mu_2 \mu_1}} \right)^{1/2}$$

Since  $\sin \theta_i \leq 1$ , we must have  $\frac{\varepsilon_2 - \mu_2}{\varepsilon_1 \mu_1} \leq \frac{\mu_1 - \mu_2}{\mu_2 \mu_1}$  or,

$$\frac{\varepsilon_2}{\varepsilon_1} \leq \frac{\mu_1}{\mu_2} \quad (\text{then there exists such a } \theta_i \text{ which makes } \Gamma_{\perp} = 0)$$

If  $\varepsilon_1 = \varepsilon_2$  and  $\mu_1 \neq \mu_2$ , then

$$\sin \theta_i = \left( \frac{1 - \frac{\mu_2}{\mu_1}}{\frac{\mu_1 - \mu_2}{\mu_2 \mu_1}} \right)^{1/2} = \left( \frac{\frac{\mu_1 - \mu_2}{\mu_1}}{\left( \frac{\mu_1^2 - \mu_2^2}{\mu_1 \mu_2} \right)} \right)^{1/2} = \left( \frac{\mu_2}{\mu_1 + \mu_2} \right)^{1/2}$$

So,

$$\sin \theta_i = \sqrt{\frac{\mu_2}{\mu_1 + \mu_2}}$$

Or,

$$\theta_i = \theta_B = \sin^{-1} \left( \sqrt{\frac{\mu_2}{\mu_1 + \mu_2}} \right) = \cos^{-1} \left( \sqrt{\frac{\mu_1}{\mu_1 + \mu_2}} \right) = \tan^{-1} \left( \sqrt{\frac{\mu_2}{\mu_1}} \right)$$

$\theta_B = \text{Brewsters Angle}$

It the two media are electrically the same but magnetically different and if the incident wave is of perpendicular polarization, then, there exists an incidence angle for which  $\Gamma_{\perp} = 0$  .

b) Parallel Polarization

$$\Gamma_T = \frac{\sqrt{\frac{\mu_2}{\epsilon_2}} \cos \theta_t - \sqrt{\frac{\mu_1}{\epsilon_1}} \cos \theta_i}{\sqrt{\frac{\mu_2}{\epsilon_2}} \cos \theta_t + \sqrt{\frac{\mu_1}{\epsilon_1}} \cos \theta_i} = 0 \quad \text{or,}$$

$$\cos \theta_i = \sqrt{\frac{\mu_2 \epsilon_1}{\mu_1 \epsilon_2}} \cos \theta_t \dots\dots\dots(1)$$

Using Snell's Law simultaneously with (1) gives:

$$\sin \theta_i = \left( \frac{\frac{\epsilon_2 - \mu_2}{\epsilon_1} - \frac{\mu_1}{\epsilon_1}}{\frac{\epsilon_2 - \epsilon_1}{\epsilon_1} - \frac{\epsilon_1}{\epsilon_1}} \right)^{1/2}$$

A real  $\theta_i$  exists if  $\frac{\epsilon_2}{\epsilon_1} - \frac{\mu_2}{\mu_1} \leq \frac{\epsilon_2}{\epsilon_1} - \frac{\epsilon_1}{\epsilon_2}$  or,

$$\frac{\mu_2}{\mu_1} \geq \frac{\epsilon_1}{\epsilon_2}$$

If  $\mu_1 = \mu_2$  and  $\varepsilon_1 \neq \varepsilon_2$  then

$$\theta_i = \theta_B = \sin^{-1} \left( \sqrt{\frac{\varepsilon_2}{\varepsilon_1 + \varepsilon_2}} \right) = \cos^{-1} \left( \sqrt{\frac{\varepsilon_1}{\varepsilon_1 + \varepsilon_2}} \right) = \tan^{-1} \left( \sqrt{\frac{\varepsilon_2}{\varepsilon_1}} \right)$$

Polarization	Medium 1	Medium 2	Does Brewster's angle exist?	$\theta_B$
Perpendicular	$\varepsilon, \mu_1$	$\varepsilon, \mu_2$	Yes	$\theta_B = \sin^{-1} \left( \sqrt{\frac{\mu_2}{\mu_1 + \mu_2}} \right) = \tan^{-1} \left( \sqrt{\frac{\mu_2}{\mu_1}} \right)$
Perpendicular	$\varepsilon_1, \mu$	$\varepsilon_2, \mu$	No	-
Parallel	$\varepsilon, \mu_1$	$\varepsilon, \mu_2$	No	-
Parallel	$\varepsilon_1, \mu$	$\varepsilon_2, \mu$	Yes	$\theta_B = \sin^{-1} \left( \sqrt{\frac{\varepsilon_2}{\varepsilon_1 + \varepsilon_2}} \right) = \tan^{-1} \left( \sqrt{\frac{\varepsilon_2}{\varepsilon_1}} \right)$